

# Vertical degree of forms and some applications

Marián Fecko

Department of Theoretical Physics

Comenius University

Bratislava

fecko@fmph.uniba.sk

Student Colloquium and School on Mathematical Physics,  
Stará Lesná, Slovakia, August 26 - September 1, 2018

## We will learn:

- What is the structure of (exterior) forms on  $L = L_1 \oplus L_2$

## We will learn:

- What is the structure of (exterior) forms on  $L = L_1 \oplus L_2$
- What is **horizontal**  $p$ -form

## We will learn:

- What is the structure of (exterior) forms on  $L = L_1 \oplus L_2$
- What is **horizontal**  $p$ -form
- What is **vertical degree**  $q$  of a  $p$ -form

## We will learn:

- What is the structure of (exterior) forms on  $L = L_1 \oplus L_2$
- What is **horizontal**  $p$ -form
- What is **vertical degree**  $q$  of a  $p$ -form
- What it looks like for **spacetime** with an **observer field**

## We will learn:

- What is the structure of (exterior) forms on  $L = L_1 \oplus L_2$
- What is **horizontal**  $p$ -form
- What is **vertical degree**  $q$  of a  $p$ -form
- What it looks like for **spacetime** with an **observer field**
- What it looks like for **principal bundle** with **connection**

## We will learn:

- What is the structure of (exterior) forms on  $L = L_1 \oplus L_2$
- What is **horizontal**  $p$ -form
- What is **vertical degree**  $q$  of a  $p$ -form
- What it looks like for **spacetime** with an **observer field**
- What it looks like for **principal bundle** with **connection**
- How it might help in computing **exterior covariant derivative**

# Contents

- 1 Introduction
- 2 Forms on  $L = L_1 \oplus L_2$
- 3 Application: Observer field  $V$  on a spacetime  $(M, g)$
- 4 Application: Connections on principal  $G$ -bundles
  - Computation of  $\Omega \equiv D\omega$
  - Computation of  $D\alpha$  for  $\alpha \in \overline{\Omega}^p(P, \rho)$



Adapted basis in  $L = L_1 \oplus L_2$ 

Let us call the subspaces  $L_1$  and  $L_2$  of  $L_1 \oplus L_2 \equiv L$

$$L \supset L_1 = \text{vertical}$$

$$L \supset L_2 = \text{horizontal}$$

Choose a basis  $e_i \in L_1$  and  $e_a \in L_2$ . Then

$$e_\alpha \equiv (e_i, e_a) = \text{adapted basis in } L = L_1 \oplus L_2$$

$$e^\alpha \equiv (e^i, e^a) = \text{adapted dual basis in } L^* = L_1^* \oplus L_2^*$$

## Duality conditions in detail

The overall duality conditions

$$\langle e^\alpha, e_\beta \rangle = \delta_\beta^\alpha \quad (1)$$

are equivalent to **four** particular cases:

$$\langle e^a, e_b \rangle = \delta_b^a \quad \langle e^i, e_j \rangle = \delta_j^i \quad \langle e^a, e_i \rangle = 0 \quad \langle e^i, e_a \rangle = 0 \quad (2)$$

Then, clearly, we can write

$$v \text{ is vertical} \Leftrightarrow v = v^i e_i \Leftrightarrow \langle e^a, v \rangle = 0 \quad (3)$$

$$u \text{ is horizontal} \Leftrightarrow u = u^a e_a \Leftrightarrow \langle e^i, u \rangle = 0 \quad (4)$$

## Forms in $L = L_1 \oplus L_2$

A **general  $p$ -form** in  $L = L_1 \oplus L_2$  reads

$$\alpha = \frac{1}{p!} \alpha_{\alpha \dots \beta} e^\alpha \wedge \dots \wedge e^\beta \quad (5)$$

( **$p$  pieces** of basis 1-forms  $e^\alpha, \dots, e^\beta$  are wedge-multiplied).

Now each  $e^\alpha$  is **either**  $e^i$  **or**  $e^a$ . So, (5) is actually a **sum** of terms of the structure

$$\alpha = \hat{\alpha} + e^i \wedge \hat{\alpha}_i + e^i \wedge e^j \wedge \hat{\alpha}_{ij} + \dots \quad (6)$$

where the hatted forms  $(\hat{\alpha}, \hat{\alpha}_i, \hat{\alpha}_{ij}, \dots)$  do not contain  $e^i$  (they only contain  $e^a$ ).

## Forms in $L = L_1 \oplus L_2$ (2)

Or, equivalently,

$$\alpha = \alpha_{(0)} + \alpha_{(1)} + \alpha_{(2)} + \dots \quad (7)$$

where

$$\alpha_{(0)} := \hat{\alpha} \quad (8)$$

$$\alpha_{(1)} := e^i \wedge \hat{\alpha}_i \quad (9)$$

$$\alpha_{(2)} := e^i \wedge e^j \wedge \hat{\alpha}_{ij} \quad (10)$$

$$\alpha_{(3)} := e^i \wedge e^j \wedge e^k \wedge \hat{\alpha}_{ijk} \quad (11)$$

$$\text{etc.} \quad (12)$$

Observation: Individual **terms** in the decomposition (7) are **well defined** (they do **not depend** on the choice of  $e_i \in L_1$  or  $e_a \in L_2$ ).

## Forms in $L = L_1 \oplus L_2$ (3)

So, a general  $p$ -form in  $L = L_1 \oplus L_2$  is a **sum** of terms characterized by **another integer** (in addition to the integer  $p$ ).

We call it vertical degree (recall that  $p$  is **the degree** of  $\alpha$ ).

Namely,

$$\text{vertical degree of } \alpha_{(q)} \text{ is (by definition) } q \in \mathbb{Z} \quad (13)$$

A general  $p$ -form in  $L = L_1 \oplus L_2$  is

(from the point of view of the **vertical** degree) **inhomogeneous**.

(It is a sum of homogeneous terms.)

## Horizontal part of a form, horizontal form

The first term of the sum is of particular interest. We call it

$$\text{horizontal part of } \alpha \text{ is (by definition) } \alpha_{(0)} \quad (14)$$

Then we introduce the concept of

$$\text{horizontal form } \alpha : \text{ such that } \alpha = \alpha_{(0)} \quad (15)$$

So, horizontal  $p$ -form is a particular homogeneous form:

$$\text{horizontal form } \alpha : \text{ such that its horizontal degree is } 0 \quad (16)$$

## Some useful observations (1)

Observation 1:

$$\alpha \text{ is horizontal} \Leftrightarrow i_{e_j} \alpha = 0 \text{ (for all } j) \quad (17)$$

That is,  $\alpha$  is **killed** by **any vertical** argument.

So, horizontal forms only can survive on horizontal arguments.

Observation 2 (more general):

$$i_{e_j} \alpha = 0 \Leftrightarrow \alpha = \alpha_{(0)} \quad (18)$$

$$i_{e_j} i_{e_k} \alpha = 0 \Leftrightarrow \alpha = \alpha_{(0)} + \alpha_{(1)} \quad (19)$$

$$i_{e_j} i_{e_k} i_{e_l} \alpha = 0 \Leftrightarrow \alpha = \alpha_{(0)} + \alpha_{(1)} + \alpha_{(2)} \quad (20)$$

$$\text{etc.} \quad (21)$$

## Some useful observations (2)

Define

$$Q := j^k i_k \quad i_k \equiv i_{e_k}, \quad j^k \alpha := e^k \wedge \alpha \quad (22)$$

Observation 3:

$$Q\alpha_{(q)} = q\alpha_{(q)} \quad (23)$$

Why?

Recall (see 5.8.11 in [1]), that

$$j^\beta i_\beta \alpha = p\alpha \quad \text{for } \alpha \text{ a } p\text{-form} \quad (24)$$

(hint:  $j^\alpha i_\beta$  is a **derivation** (of the algebra of forms) of **degree 0**).For  $j^k i_k$  similarly. (It works like **number operator**  $a_k^+ a_k$  :-)



## Observer field $V$

Let  $(M, g)$  be a **space-time** (signature  $+ - - -$ ) and let  $V$  be a future-oriented vector field with  $g(V, V) = 1$ . Integral curves of  $V$  may be regarded as **world-lines** of **observers**.

In each (world-)point,

- $V \equiv e_0$  defines (observer's, local) **time** direction
- $V^\perp$  defines (observer's, local) **3-space**.

So, we can regard (in **tangent** space)

$$L_1 := \text{Span } V \quad L_2 := V^\perp \equiv \text{the 3-space} \quad (25)$$

Then

$$e_i \leftrightarrow V \equiv e_0 \quad e^i \leftrightarrow \tilde{V} \equiv g(V, \cdot) \equiv e^0 \quad (26)$$

## Observer field $V$ and 1+3 decomposition

Since there is **just a single**  $e^i \leftrightarrow \tilde{V}$ , the general expression

$$\alpha = \alpha_{(0)} + \alpha_{(1)} + \alpha_{(2)} + \dots \quad (27)$$

$$= \hat{\alpha} + e^i \wedge \hat{\alpha}_i + e^i \wedge e^j \wedge \hat{\alpha}_{ij} + \dots \quad (28)$$

reduces, here, to just **two terms**

$$\alpha = \alpha_{(0)} + \alpha_{(1)} \quad (29)$$

$$= \hat{r} + \tilde{V} \wedge \hat{s} \quad (30)$$

A general form is given (w.r.t.  $V$ )  
by **a pair of spatial** (= horizontal) forms  $(\hat{s}, \hat{r})$ .

# Principal $G$ -bundle - adapted frame and co-frame fields (1)

Let

- $G$  be a semi-simple Lie group,  $\mathcal{G}$  its Lie algebra
- $E_i$  an orthonormal basis in  $\mathcal{G}$  (w.r.t. Killing-Cartan  $K$ )
- $\pi : P \rightarrow M$  be a **principal  $G$ -bundle** over space-time  $(M, \hat{g})$
- $\omega = \omega^i E_i$  be a **connection form**
- $\xi_X, X = X^i E_i$ , be the **fundamental field** of  $R_g : P \rightarrow P$
- $\xi_{E_i}$  be the **generators** of  $R_g : P \rightarrow P$
- $w^h$  be the **horizontal lift** of a vector field  $w$  on  $M$
- $\hat{e}_a$  be an orthonormal (w.r.t.  $\hat{g}$ ) frame on (a part of)  $M$

## Principal $G$ -bundle - adapted frame and co-frame fields (2)

Then

$$e_\alpha \equiv (e_i, e_a) \quad e^\alpha \equiv (e^i, e^a) \quad \langle e^\alpha, e_\beta \rangle = \delta_\beta^\alpha$$

where

$$e_i := \xi E_i \quad e_a := \hat{e}_a^h \quad (31)$$

$$e^i := \omega^i \quad e^a := \pi^* \hat{e}^a \quad (32)$$

constitute a **frame** and the **dual co-frame**, respectively (on part of  $P$ ).

Moreover, the frames happen to be **ortho-normal** w.r.t.

$$g := \pi^* \hat{g} + K(\omega, \omega) \quad (33)$$

## Principal $G$ -bundle - adapted frame and co-frame fields (3)

Now, in each tangent space  $L \equiv T_p P$  in point  $p \in P$ , define

$$L_1 := \text{ver } T_p P \quad L_2 := \text{hor } T_p P \quad (34)$$

in the sense of **connection theory** (!).

Then the decomposition (the essence of **connection** concept)

$$T_p P = \text{ver } T_p P \oplus \text{hor } T_p P \quad (35)$$

**exactly matches** the decomposition studied above

$$L = L_1 \oplus L_2 \quad (36)$$

(including **meaning** of hor and ver).

# Principal $G$ -bundle - adapted frame and co-frame fields (4)

So, as an example,

- the **curvature** 2-forms  $\Omega^i = \frac{1}{2}\Omega^i_{ab}e^a \wedge e^b$  are **horizontal** ( $\equiv \alpha_{(0)}$ )
- the **connection** 1-forms  $\omega^i$  have **vertical degree equal 1** ( $\equiv \alpha_{(1)}$ )
- the expression  $\frac{1}{2}c^i_{jk}\omega^j \wedge \omega^k$  has **vertical degree equal 2** ( $\equiv \alpha_{(2)}$ )
- and so on

# Computation of $\Omega \equiv D\omega$

By definition, the **curvature** 2-form  $\Omega$  is the **exterior covariant** derivative

$$D := \text{hor} \circ d \quad (37)$$

of the **connection form**  $\omega \equiv \omega^i E_i$ .

So, we need to compute

$$\Omega^i := \text{hor}(d\omega^i) \equiv (de^i)_{(0)} \quad (38)$$

# Computation of $\Omega \equiv D\omega$ (2)

**Definition** properties of connection form  $\omega$ :

$$R_g^* \omega = \text{Ad}_{g^{-1}} \omega \quad (39)$$

$$\langle \omega, \xi_X \rangle = X \quad (40)$$

Their **infinitesimal** version

$$\mathcal{L}_{\xi_X} \omega = -\text{ad}_X \omega \equiv -[X, \omega] \quad (41)$$

$$i_{\xi_X} \omega = X \quad (42)$$

is **equivalent** to

$$i_{\xi_X} d\omega = -[X, \omega] \quad (43)$$

$$i_{\xi_X} \omega = X \quad (44)$$

(use  $\mathcal{L}_W = di_W + i_W d$  and (42))



# Computation of $\Omega \equiv D\omega$ (3)

Finally (for  $X = E_j$ ,  $\omega = \omega^i E_i$ ,  $e^i = \omega^i$ ) we get

$$i_{e_j}(de^i) = -c_{jk}^i e^k \quad (45)$$

$$i_{e_j} e^i = \delta_j^i \quad (46)$$

Consequently,

$$i_{e_k} i_{e_j}(de^i) = -c_{jk}^i$$

$$i_{e_m} i_{e_k} i_{e_j}(de^i) = 0$$

(47)

## Computation of $\Omega \equiv D\omega$ (4)

From the second formula in (47) we can, first, deduce

$$de^i = (de^i)_{(0)} + (de^i)_{(1)} + (de^i)_{(2)} \quad (48)$$

From the first formula in (47), then, we have

$$(de^i)_{(2)} = -\frac{1}{2}c_{jk}^i e^j \wedge e^k \quad (49)$$

and, finally, from (45),

$$(de^i)_{(1)} = 0 \quad (50)$$

(no **vertical degree one** term is at the r.h.s of (45)).

## Computation of $\Omega \equiv D\omega$ (5)

So, we learned that

$$de^i = (de^i)_{(0)} - \frac{1}{2}c_{jk}^i e^j \wedge e^k \quad (51)$$

or, already trivially,

$$(de^i)_{(0)} =: \text{hor } de^i = de^i + \frac{1}{2}c_{jk}^i e^j \wedge e^k \quad (52)$$

So, remembering that  $e^i = \omega^i$ , we get well-known explicit formula

$$\Omega^i := \text{hor } d\omega^i = d\omega^i + \frac{1}{2}c_{jk}^i \omega^j \wedge \omega^k \quad (53)$$

## Bianchi identity: $D\Omega = 0$

In our language:

$$DDe^i = 0 \quad \text{i.e. } (dDe^i)_{(0)} = 0 \quad (54)$$

Why is it true?

Because actually

$$dDe^i = (dDe^i)_{(1)} + (dDe^i)_{(3)} \quad (55)$$

Indeed,

$$dDe^i = d(de^i + \frac{1}{2}c_{jk}^i e^j \wedge e^k) = c_{jk}^i (de^j) \wedge e^k \quad (56)$$

Combine with (68)  $\Rightarrow$  :-)

# Computation of $D\alpha$ (1)

**Definition** properties of a **horizontal  $p$ -form  $\alpha$  of type  $\rho$** :

$$R_g^* \alpha = \rho(g^{-1})\alpha \quad (\text{type } \rho) \quad (57)$$

$$i_{\xi_X} \alpha = 0 \quad (\text{horizontal}) \quad (58)$$

Their **infinitesimal** version

$$\mathcal{L}_{\xi_X} \alpha = -\rho'(X)\alpha \quad (59)$$

$$i_{\xi_X} \alpha = 0 \quad (60)$$

is **equivalent** to

$$i_{\xi_X}(d\alpha) = -\rho'(X)\alpha \quad (61)$$

$$i_{\xi_X} \alpha = 0 \quad (62)$$

## Computation of $D\alpha$ (2)

Finally (for  $X = E_j$ ) we get

$$i_{E_j}(d\alpha) = -\rho'(E_j)\alpha \quad (63)$$

$$i_{E_j}\alpha = 0 \quad (64)$$

Consequently,

$$i_{E_k}i_{E_j}(d\alpha) = 0 \quad (65)$$

From (65) we, first, deduce

$$d\alpha = (d\alpha)_{(0)} + (d\alpha)_{(1)} \quad (66)$$

and from (63) we see that

$$(d\alpha)_{(1)} = -e^i \wedge \rho'(E_i)\alpha \quad (67)$$

## Computation of $D\alpha$ (3)

Altogether, we learned that

$$d\alpha = (d\alpha)_{(0)} - e^i \wedge \rho'(E_i)\alpha \quad (68)$$

or, trivially,

$$(d\alpha)_{(0)} =: \text{hor } d\alpha = d\alpha + e^i \wedge \rho'(E_i)\alpha \quad (69)$$

So, remembering that  $e^i = \omega^i$ , we get well-known explicit formula

$$D\alpha := \text{hor } d\alpha = d\alpha + \omega^i \wedge \rho'(E_i)\alpha \equiv d\alpha + \rho'(\omega)\dot{\wedge}\alpha \quad (70)$$

# References (1)



M. Fecko.

*Differential Geometry and Lie Groups for Physicists.*

Cambridge University Press, Cambridge, UK (2006, 2011)



# References (1)



M. Fecko.

*Differential Geometry and Lie Groups for Physicists.*

Cambridge University Press, Cambridge, UK (2006, 2011)



M. Dubois-Violette, J. Madore

*Conservation Laws and Integrability Conditions for  
Gravitational and Yang-Mills Field Equations.*

Commun. Math. Phys. 108, (1987) 213-223

# References (1)



M. Fecko.

*Differential Geometry and Lie Groups for Physicists.*

Cambridge University Press, Cambridge, UK (2006, 2011)



M. Dubois-Violette, J. Madore

*Conservation Laws and Integrability Conditions for  
Gravitational and Yang-Mills Field Equations.*

Commun. Math. Phys. 108, (1987) 213-223