

FUZZY ONION STRIKES BACK

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Take home message



TAKE HOME MESSAGE

- Quantization of gravity seems to lead to quantized spacetimes.
- Fuzzy onion is a three dimensional model of such a space.
- It is rather straightforward to work with so it is a nice toy model / playground to check the consequences of quantum structure.



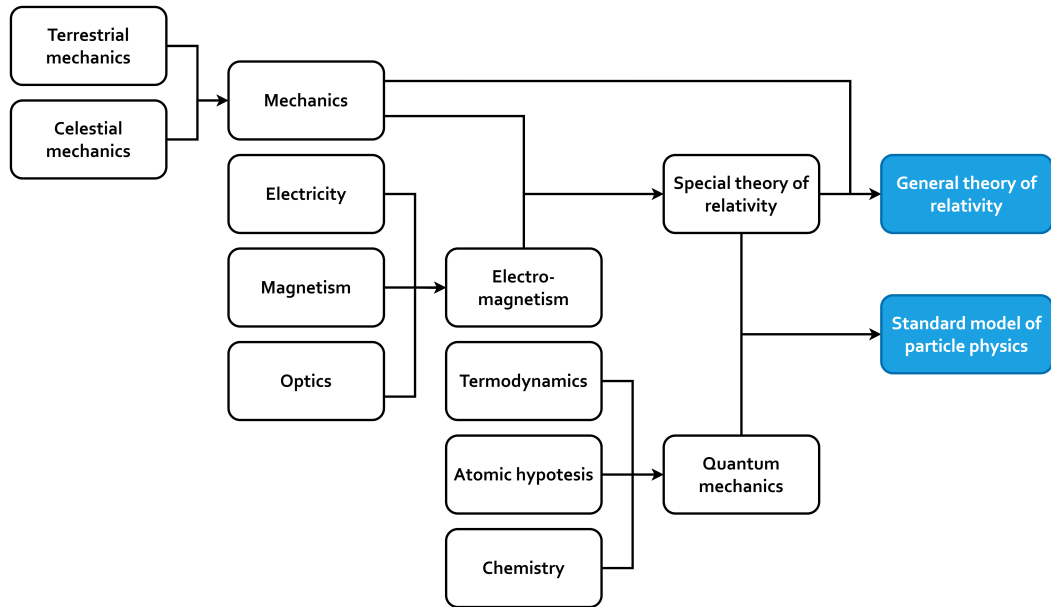
TAKE HOME MESSAGE

- Quantization of gravity seems to lead to quantized spacetimes.
- Fuzzy onion is a three dimensional model of such a space.
- It is rather straightforward to work with so it is a nice toy model / playground to check the consequences of quantum structure.
- We do have PhD. and postdoc positions related to this and more in Bratislava.



Quick motivation





QUANTUM STRUCTURE OF SPACETIME

- We need a quantum theory of gravity.
- Quantization of general relativity leads to a nonrenormalizable theory.
- We have reasons to believe that future theory of quantum gravity will have a different notion of spacetime.
No distinction between points under certain length scales. [[Hossenfelder 1203.6191](#)]
- Reasons:
 - gravitational Heisenberg microscope,
 - instability of quantum gravitational vacuum, [[Doplicher, Fredenhagen, Roberts '95](#)]
 - emergent spacetime.

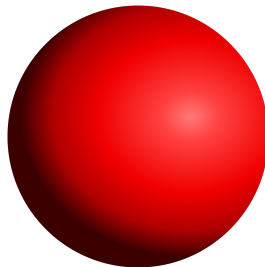
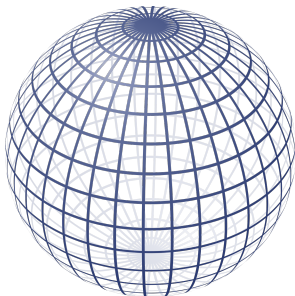


Fuzzy spaces



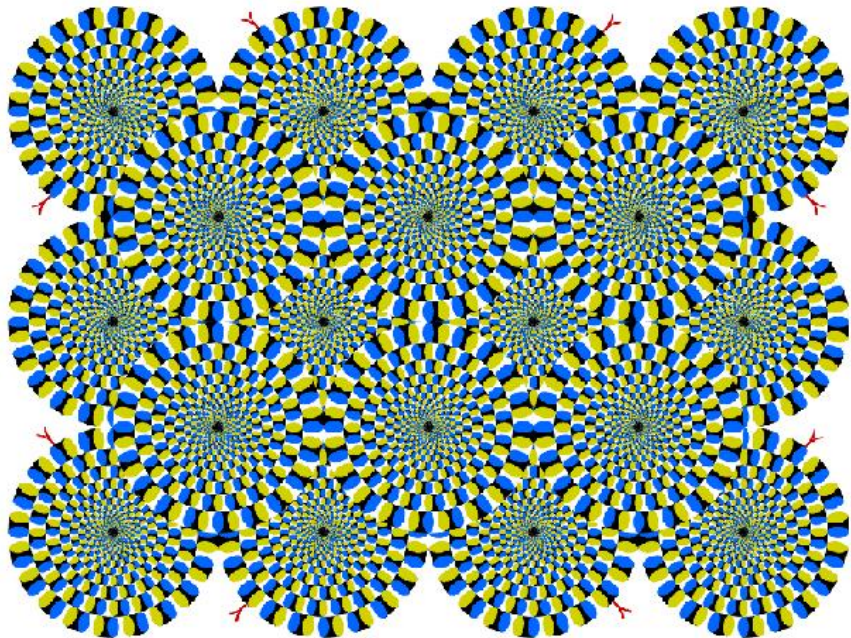
FUZZY SPHERE

- The sphere is divided into N cells. Function on the fuzzy sphere is given by a matrix Φ and the eigenvalues of Φ represent the values of the function on these cells.



- However there are no sharp boundaries between the pieces and everything is blurred, or fuzzy.





- The regular sphere S^2 is given by the coordinates

$$x_i x_i = R^2 \quad , \quad x_i x_j - x_j x_i = 0 \quad , \quad i, j = 1, 2, 3 \quad ,$$

which generate the algebra of functions.

- For the fuzzy sphere S_N^2 we define

$$\hat{x}_i \hat{x}_i = r^2 \quad , \quad \hat{x}_i \hat{x}_j - \hat{x}_j \hat{x}_i = i\theta \varepsilon_{ijk} \hat{x}_k \quad , \quad i, j = 1, 2, 3 \quad .$$

- Such \hat{x}_i 's generate a different, non-commutative, algebra and S_N^2 is an object, which has this algebra as an algebra of functions.



- The conditions can be realized as an $N = 2s + 1$ dimensional representation of $SU(2)$

$$\hat{x}_i = \frac{2r}{\sqrt{N^2 - 1}} L_i \quad , \quad \theta = \frac{2r}{\sqrt{N^2 - 1}} \sim \frac{2}{N} \quad , \quad \rho^2 = \frac{4r^2}{N^2 - 1} s(s + 1) = r^2 \quad .$$

- The group $SU(2)$ still acts on \hat{x}_i 's and this space enjoys a full rotational symmetry.
- (Matrix) geometry encoded in the Laplacian

$$\mathcal{K}\Phi = \frac{1}{r^2} [L_i, [L_i, \Phi]] \quad .$$

Analogue of the integral is the trace

$$\frac{4\pi r^2}{N} \text{tr}_N(\Phi) \quad .$$

- In the limit $N \rightarrow \infty$ we recover the original sphere.



- Most importantly nonzero commutators

$$\hat{x}_i \hat{x}_i = \rho^2 \quad , \quad \hat{x}_i \hat{x}_j - \hat{x}_j \hat{x}_i = i\theta \varepsilon_{ijk} \hat{x}_k \quad , \quad i = 1, 2, 3 \quad .$$

imply uncertainty relations for positions

$$\Delta x_i \Delta x_j \neq 0 \quad .$$

- Configuration space is analogous to phase space of quantum mechanics.
- In a similar fashion it is possible to construct an analogous deformation of the plane

$$\hat{x}_i \hat{x}_j - \hat{x}_j \hat{x}_i = i\theta \varepsilon_{ij} = i\theta_{ij} \quad , \quad i = 1, 2 \quad .$$

Construction uses the \star -product

$$f \star g = f e^{\frac{i}{2} \overleftarrow{\partial} \theta \overrightarrow{\partial}} g = fg + \frac{i\theta^{\mu\nu}}{2} \frac{\partial f}{\partial x^\mu} \frac{\partial g}{\partial x^\nu} + \dots$$



- Regularization of infinities in the standard QFT.
[\[Heisenberg ~1930; Snyder 1947, Yang 1947\]](#)
- Regularization of field theories for numerical simulations.
[\[Panero 2016\]](#)
- An effective description of the open string dynamics in a magnetic background in the low energy limit.
[\[Seiberg Witten 1999; Douglas, Nekrasov 2001\]](#)
- Solutions of various matrix formulations of the string theory (IKKT, BFSS, BNM).
[\[Steinacker 2013\]](#)
- Geometric unification of the particle physics and theory of gravity.
[\[van Suijlekom 2015\]](#)
- An effective description of various systems in a certain limit (eg. QHE).
[\[Karabali, Nair 2006\]](#)
- Toy models of spaces with discrete quantum structure, which is expected to arise in quantum theory of gravity.





The/A fuzzy onion



THE/A FUZZY ONION

- Idea: construct a 3D non-commutative space as a series of concentric fuzzy spheres of increasing radius.



[Hammou, Lagraa, Sheikh-Jabbari 2002; Vitale, Wallet 2013; Scholtz et. al; Schupp, Solodukhin 2009; Prešnajder, Gáliková, Kováčik 2015]



THE/A FUZZY ONION

Construction of [Prešnajder, Gáliková, Kováčik 2015; Scholtz et. al].

- Start with

$$[x_i, x_j] = 2\lambda i \varepsilon_{ijk} x_k$$

- Explicitly realize by

$$\hat{x}_i = \frac{\theta}{2} a_\alpha^\dagger \sigma_{\alpha\beta}^i a_\beta .$$

where

$$[a_\alpha, a_\beta^\dagger] = \delta_{\alpha\beta}, \quad [a_\alpha, a_\beta] = [a_\alpha^\dagger, a_\beta^\dagger] = 0 .$$

- Free Hamiltonian, i.e. the Laplacian

$$H_0 \Psi = \frac{1}{2\lambda r} [a_\alpha^\dagger, [a_\alpha, \Psi]] , \quad r = \lambda (a_\alpha^\dagger a_\alpha + 1) .$$

- Spectrum of hydrogen atom

$$E_{\lambda n}^I = \frac{\hbar}{m_e \lambda^2} \left(1 - \sqrt{1 + \left(\frac{m_e q \lambda}{\hbar^2 n} \right)^2} \right) .$$



THE/A FUZZY ONION

- Ours is a bottom-up approach.
- Take M fuzzy spheres or radii $r = \lambda, 2\lambda, \dots, M\lambda$.
- Functions given by a matrix

$$\Psi = \begin{pmatrix} \Phi^{(1)} & & & \\ & \Phi^{(2)} & & \\ & & \ddots & \\ & & & \Phi^{(M)} \end{pmatrix} .$$

- Recall the single layer expression

$$\theta = \frac{2r}{\sqrt{N^2 - 1}} .$$

- The dimension of this space is

$$d = \sum_{N=1}^M N^2 = \frac{1}{6} M(M+1)(2M+1) .$$



THE/A FUZZY ONION

- Recall the standard three-dimensional integration of a function ψ

$$\int d^3x \, \psi = \int r^2 \, dr \int d\Omega \, \psi .$$

- We change this to a version discrete in the radial direction

$$\sum_{N=1}^M (\lambda N)^2 \lambda \frac{4\pi}{N} \operatorname{tr}_N \Phi^{(N)} = \operatorname{Tr} (4\pi \lambda^2 r \, \Psi)$$

where we have defined the radial distance matrix r as

$$r = \begin{pmatrix} \lambda \, \mathbb{1}_{1 \times 1} & & & & \\ & 2\lambda \, \mathbb{1}_{2 \times 2} & & & \\ & & 3\lambda \, \mathbb{1}_{3 \times 3} & & \\ & & & \ddots & \\ & & & & M\lambda \, \mathbb{1}_{M \times M} \end{pmatrix}$$



- The angular part of the kinetic term defined layerwise

$$\mathcal{K}_L \Psi = r^{-2} \begin{pmatrix} \mathcal{K}^{(1)} \Phi^{(1)} & & & \\ & \mathcal{K}^{(2)} \Phi^{(2)} & & \\ & & \mathcal{K}^{(3)} \Phi^{(3)} & \\ & & & \ddots \\ & & & & \mathcal{K}^{(M)} \Phi^{(M)} \end{pmatrix} .$$

- What about the radial direction?



THE/A FUZZY ONION

$$\text{for } \Phi^{(N)} = \sum_{l=0}^{N-1} \sum_{m=-l}^l c_{lm}^{(N)} Y_{lm}^{(N)}, \quad \Phi^{(N+1)} = \sum_{l=0}^N \sum_{m=-l}^l c_{lm}^{(N+1)} Y_{lm}^{(N+1)}$$

$$\mathcal{D} : \Phi^{(N+1)} \rightarrow \Phi^{(N)} = \sum_{l=0}^{N-1} \sum_{m=-l}^l c_{lm}^{(N)} Y_{lm}^{(N)}, \quad c_{lm}^{(N)} = c_{lm}^{(N+1)} \text{ for } l \leq N-1$$

$$\mathcal{U} : \Phi^{(N)} \rightarrow \Phi^{(N+1)} = \sum_{l=0}^N \sum_{m=-l}^l c_{lm}^{(N+1)} Y_{lm}^{(N+1)}, \quad \left\{ \begin{array}{l} c_{lm}^{(N+1)} = c_{lm}^{(N)} \text{ for } l \leq N-1 \\ c_{Nm}^{(N+1)} = 0 \end{array} \right. .$$



THE/A FUZZY ONION

- Define the first and second radial derivatives

$$\partial_r^{(N)} \Phi^{(N)} = \frac{\mathcal{D}\Phi^{(N+1)} - \mathcal{U}\Phi^{(N-1)}}{2\lambda},$$

$$\partial_r^2 \Phi^{(N)} = \frac{\mathcal{D}\Phi^{(N+1)} - 2\Phi^{(N)} + \mathcal{U}\Phi^{(N-1)}}{\lambda^2}.$$

- Define the radial part of Laplacian

$$\mathcal{K}_R \Psi = \partial_r^2 \Psi + 2r^{-1} \partial_r \Psi, \quad \partial_r \Psi = \begin{pmatrix} \partial_r^{(1)} \Phi^{(1)} & & & & \\ & \partial_r^{(2)} \Phi^{(2)} & & & \\ & & \partial_r^{(3)} \Phi^{(3)} & & \\ & & & \ddots & \\ & & & & \partial_r^{(M)} \Phi^{(M)} \end{pmatrix}.$$

- Recall

$$\frac{f(x+\varepsilon) - f(x-\varepsilon)}{2\varepsilon} \rightarrow f'(x), \quad \frac{f(x+\varepsilon) - 2f(x) + f(x-\varepsilon))}{\varepsilon^2} \rightarrow f''(x), \quad \Delta = r^{-2} \partial_r r^2 \partial_r + \Delta_\Omega.$$



THE/A FUZZY ONION

- Define the first and second radial derivatives

$$\begin{aligned}\partial_r^{(N)}\Phi^{(N)} &= \frac{\mathcal{D}\Phi^{(N+1)} - \mathcal{U}\Phi^{(N-1)}}{2\lambda} , \\ \partial_r^2{}^{(N)}\Phi^{(N)} &= \frac{\mathcal{D}\Phi^{(N+1)} - 2\Phi^{(N)} + \mathcal{U}\Phi^{(N-1)}}{\lambda^2} .\end{aligned}$$

- Define the radial part of Laplacian

$$\mathcal{K}_R\Psi = \partial_r^2\Psi + 2r^{-1}\partial_r\Psi , \quad \partial_r\Psi = \begin{pmatrix} \partial_r^{(1)}\Phi^{(1)} & & & & \\ & \partial_r^{(2)}\Phi^{(2)} & & & \\ & & \partial_r^{(3)}\Phi^{(3)} & & \\ & & & \ddots & \\ & & & & \partial_r^{(M)}\Phi^{(M)} \end{pmatrix} .$$

- Laplace operator for the/a fuzzy onion

$$\mathcal{K} = \mathcal{K}_R + \mathcal{K}_L .$$



Some onion physics



Quantum mechanical problems



- Hamiltonian

$$H = -\frac{\hbar^2}{2m_e}\mathcal{K} - qr^{-1}$$

with $\hbar = m_e = q = 1$, i.e. eigenvalue problem

$$H\Psi = E\Psi .$$

- We can express H as a matrix acting on vectors

$$\mathcal{C}^T = \left(c_{00}^{(1)}, c_{00}^{(2)}, c_{1-1}^{(2)}, c_{10}^{(2)}, c_{11}^{(2)}, \dots \right)^T$$

with c 's from the decomposition

$$\Phi^{(N)} = \sum_{l=0}^{N-1} \sum_{m=-l}^l c_{lm}^{(N)} Y_{lm}^{(N)} .$$



- Hamiltonian is now expressed as a $d \times d$ matrix

$$H = -\frac{1}{2}K - r^{-1} .$$

- We can make restriction to $l = 0$ sector thanks to rotational symmetry to make it $M \times M$.
- Look for eigenvalues of H .



QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$M = 50$ and $\lambda = 1$

n	1	2	3	4	5	6
E_n	-0.4142	-0.1180	-0.0541	-0.0307	-0.0179	-0.0031
$E_{\lambda n}^I$	-0.4142	-0.1180	-0.0541	-0.0307	-0.0198	-0.0138
E_n^{CQM}	-0.5	-0.125	-0.0556	-0.0313	-0.02	-0.0139



QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$$n = 1$$

M	$\lambda = 0.1$	$\lambda = 0.01$	$\lambda = 0.001$
50	$6.24 \cdot 10^{-3}$	N/A	N/A
100	$1.27 \cdot 10^{-6}$	N/A	N/A
200	$1.97 \cdot 10^{-13}$	2.81	N/A
400	$1.56 \cdot 10^{-13}$	$3.41 \cdot 10^{-2}$	N/A
800	$5.22 \cdot 10^{-13}$	$4.9 \cdot 10^{-5}$	N/A
1600	$4.8 \cdot 10^{-14}$	$1.13 \cdot 10^{-11}$	N/A
3200	$9.02 \cdot 10^{-15}$	$5.75 \cdot 10^{-12}$	$1.26 \cdot 10^{-1}$

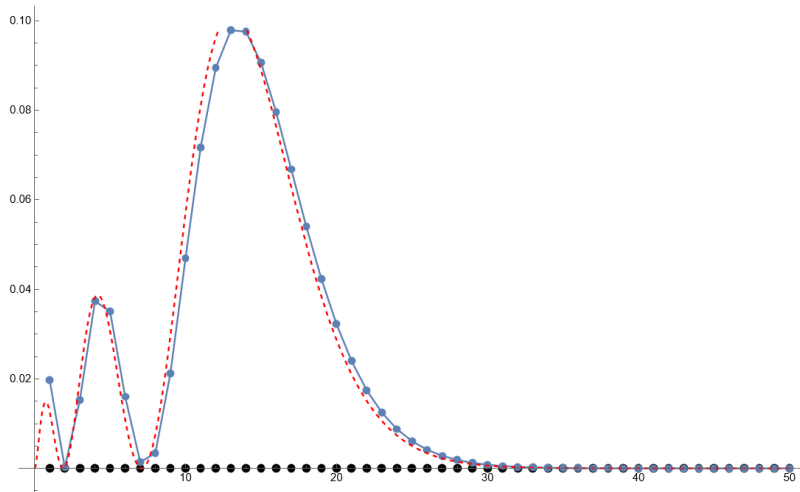


- It seems that in the limit $M \rightarrow \infty$ we recover the construction of [\[Prešnajder, Gáliková, Kováčik 2015\]](#).
- We can go further and look at the wavefunctions.



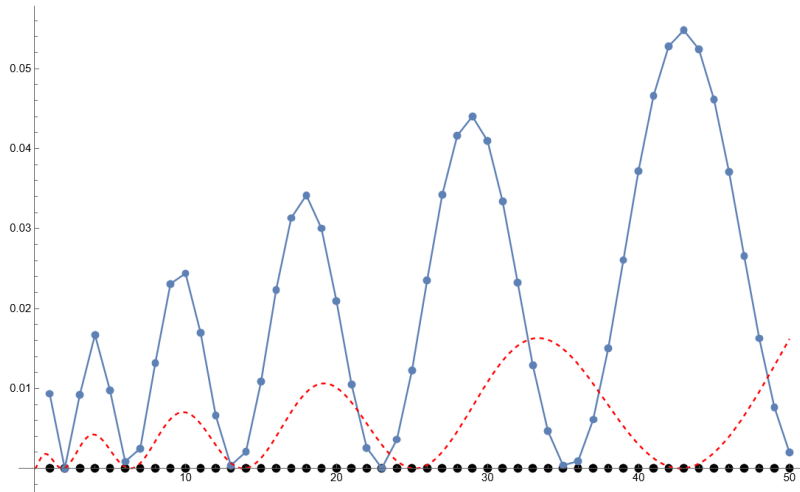
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$M = 50, \lambda = 1, n = 3$



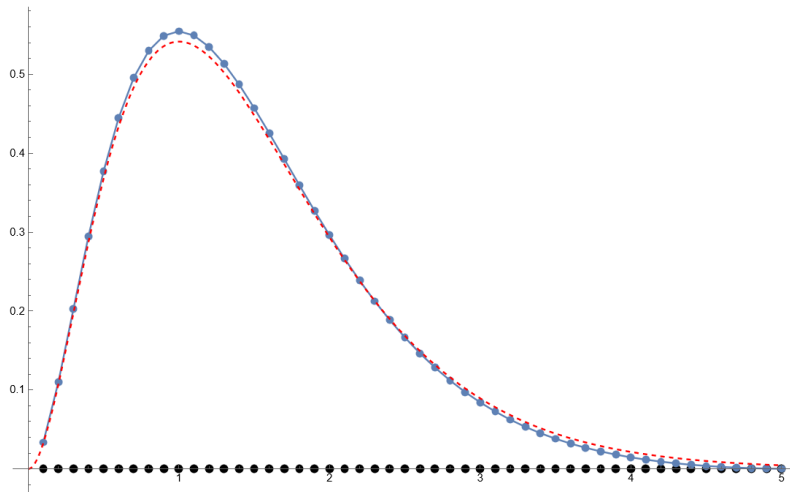
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$$M = 50, \lambda = 1, n = 6$$



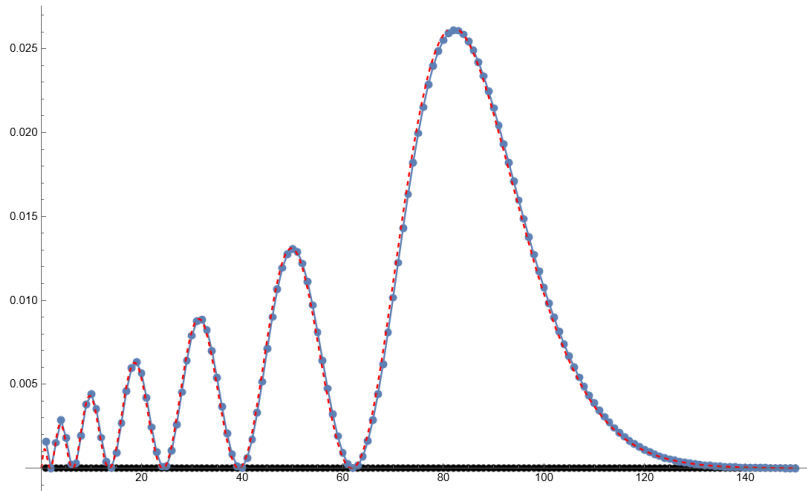
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$M = 50, \lambda = 0.1, n = 1$



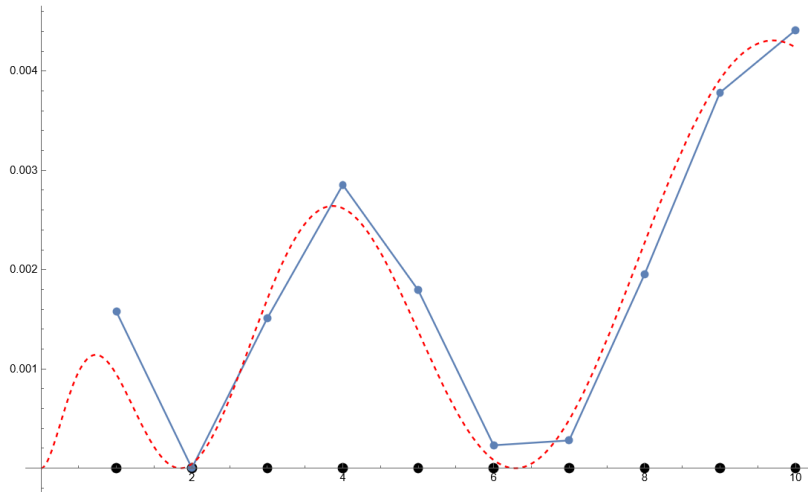
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$$M = 150, \lambda = 1, n = 7$$



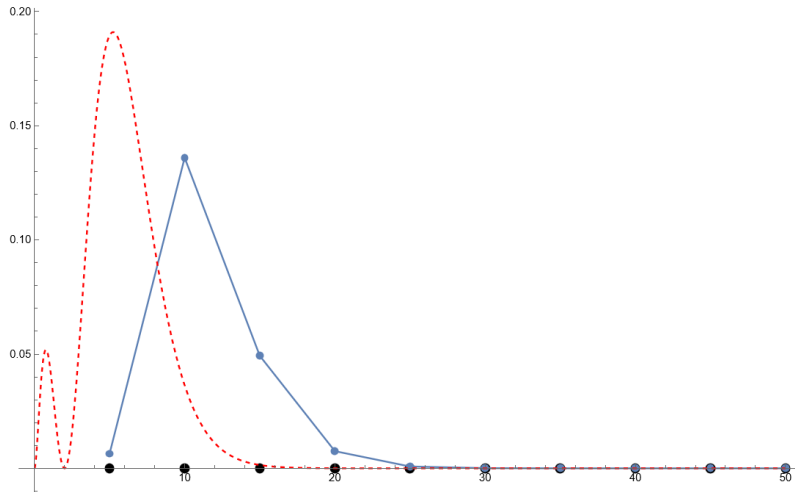
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$M = 150, \lambda = 1, n = 7$



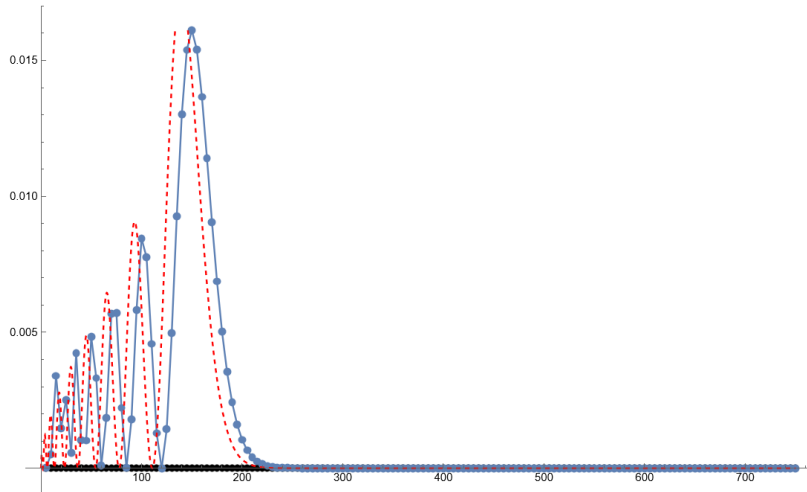
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$$M = 50, \lambda = 5, n = 2$$



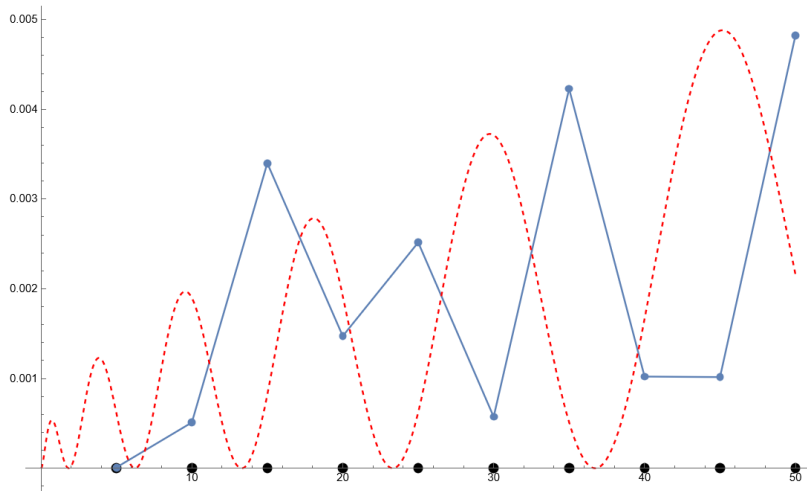
QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$$M = 150, \lambda = 5, n = 9$$



QUANTUM MECHANICAL PROBLEMS – HYDROGEN ATOM

$M = 150, \lambda = 5, n = 9$



QUANTUM MECHANICAL PROBLEMS – HARMONIC OSCILLATOR

- The same Hamiltonian with

$$H = -\frac{1}{2}K + \frac{1}{2}r^2$$

describes three dimensional harmonic oscillator.

- This can not be solved analytically in the previous approach. $M = 100$, $\lambda = 1$, $l = 0, 1$

n	0	1	2	3	4	5
E_n	1.4984	3.4922	5.4809	7.4645	9.4431	11.4166
E_n^{CQM}	1.5	3.5	5.5	7.5	9.5	11.5

n	0	1	2	3	4	5
E_n	2.5005	4.4979	6.491	8.4795	10.4632	12.4422
E_n^{CQM}	2.5	4.5	6.5	8.5	10.5	12.5



QUANTUM MECHANICAL PROBLEMS

- In principle any potential can be analyzed exactly.
- The only problem is how to recover the limit of classical (and infinite) space.



- Fuzzy onion is a regularization of a spherical cavity. In the limit $M \rightarrow \infty, \lambda \rightarrow 0, M\lambda \rightarrow R$ we recover continuous cavity of radius R .
- In chemical literature this models atoms under pressure. [\[refs in \[2\]\]](#)



QUANTUM MECHANICAL PROBLEMS – FUZZY CAVITY

S orbital [2]

$r_0 \backslash M$	∞ as a reference	100	1000	10000
0.5	14.747970030350280	14.406037740780091	14.713401904425357	14.744509454556344
1	2.373990866103664	2.300565723232022	2.366554263053759	2.373246264259394
3	-0.423967287733454	-0.427225951376656	-0.424313148630359	-0.424002075109953
10	-0.499999263281525	-0.498755577647694	-0.499986776756742	-0.499999139626354
20	-0.4999999999999994	-0.495097567963923	-0.499950009998093	-0.499999499991737



QUANTUM MECHANICAL PROBLEMS – FUZZY CAVITY

P orbitals [2]

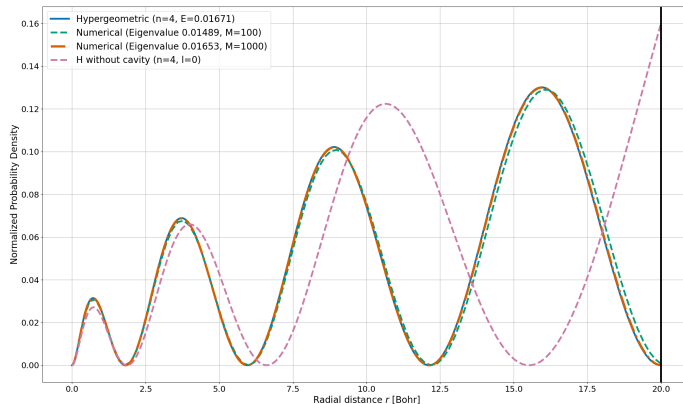
$r_0 \backslash M$	∞ as a reference	100	1000	10000
0.5	72.672039190463577	71.154357099658682	72.520341020495181	72.656870688127995
1	16.570256093469736	16.206670008470599	16.533894253548414	16.566620019574611
3	1.111684737436364	1.078613638687640	1.108361257317863	1.111352239065044
10	-0.112806210295841	-0.113415153701996	-0.112878188197422	-0.112813520180460
20	-0.124987114312918	-0.124677720985899	-0.124984183578311	-0.124987102906836

$r_0 \backslash M$	∞ as a reference	100	1000	10000
0.5	36.658875880189399	35.160313617726310	36.505181049931707	36.643467765541793
1	8.223138316160864	7.866678332336676	8.186560245882733	8.219471198758979
3	0.481250312526643	0.449669060926531	0.478000341365532	0.480924423219991
10	-0.118859544853860	-0.119527885630338	-0.118934647870327	-0.118867073891238
20	-0.124994606647078	-0.124692831261815	-0.124995259742673	-0.124994633404707



QUANTUM MECHANICAL PROBLEMS – FUZZY CAVITY

[2]



Scalar field theory



- Defined by the (euclidean) action

$$S[\Psi] = 4\pi\lambda^2 \text{Tr} [r (a \Psi \mathcal{K} \Psi + b \Psi^2 + c \Psi^4)]$$

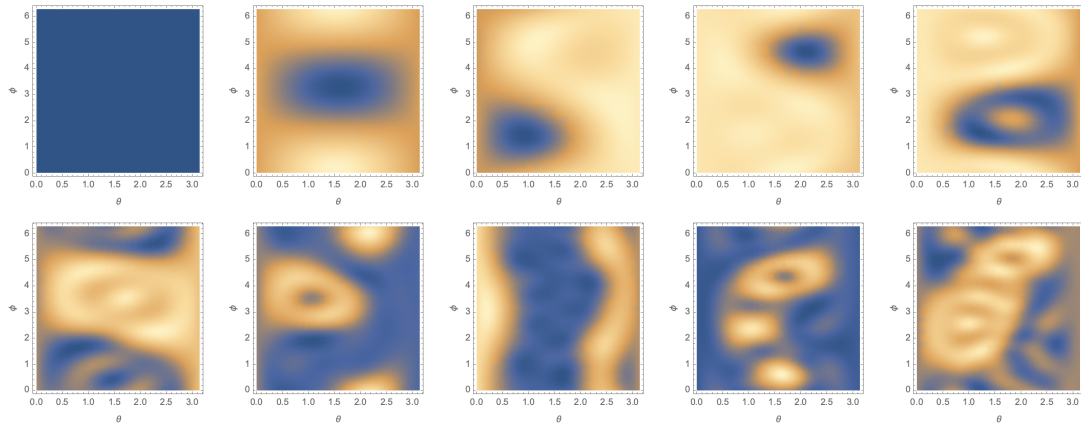
and expectation values

$$\langle \mathcal{O}(\Psi) \rangle = \frac{1}{Z} \int d\Psi e^{-S(\Psi)} \mathcal{O}(\Psi) , \quad d\Psi = \prod_{N=1}^M d\Phi^{(N)} .$$

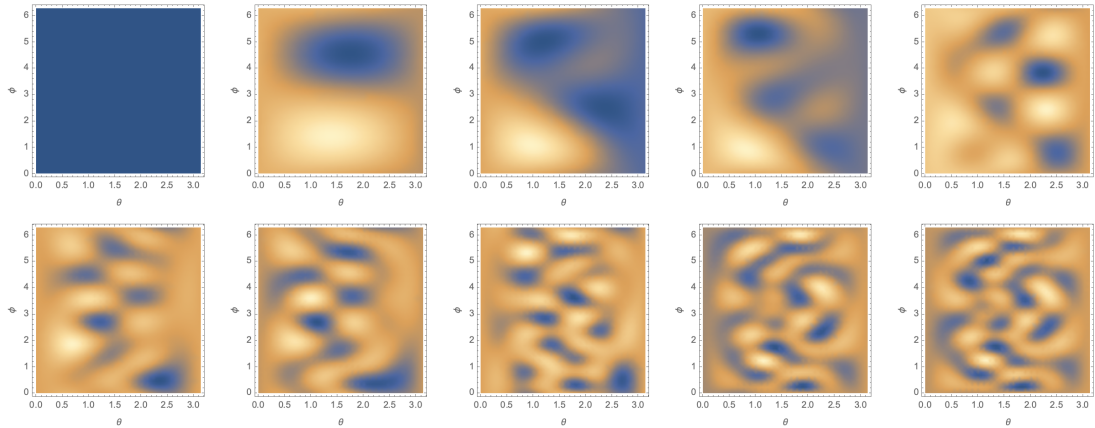
- Hybrid Monte Carlo evolution of a field configuration.[\[1\]](#)



SCALAR FIELD THEORY



SCALAR FIELD THEORY



- Radial part of the Laplacian couples the oscillations in also in radial direction.



- Radial part of the Laplacian couples the oscillations in also in radial direction.
- What about symmetry breaking? How – if at all – do the phases on layers align? Is the derivative enough or do we need something further?



Different Laplacians



- Recall that the definition of Laplacian using

$$\partial_r^{(N)} \Phi^{(N)} = \frac{\mathcal{D}\Phi^{(N+1)} - \mathcal{U}\Phi^{(N-1)}}{2\lambda},$$
$$\partial_r^{2(N)} \Phi^{(N)} = \frac{\mathcal{D}\Phi^{(N+1)} - 2\Phi^{(N)} + \mathcal{U}\Phi^{(N-1)}}{\lambda^2},$$

and

$$\mathcal{K}_R \Psi = \partial_r^2 \Psi + 2r^{-1} \partial_r \Psi$$

was in some sense arbitrary.



DIFFERENT LAPLACIANS

- Recall that the definition of Laplacian using

$$\partial_r^{(N)} \Phi^{(N)} = \frac{\mathcal{D}\Phi^{(N+1)} - \mathcal{U}\Phi^{(N-1)}}{2\lambda},$$
$$\partial_r^{2(N)} \Phi^{(N)} = \frac{\mathcal{D}\Phi^{(N+1)} - 2\Phi^{(N)} + \mathcal{U}\Phi^{(N-1)}}{\lambda^2},$$

and

$$\mathcal{K}_R \Psi = \partial_r^2 \Psi + 2r^{-1} \partial_r \Psi$$

was in some sense arbitrary.

- What are consequences of other choices? Any preferred choice?

$$\mathcal{K}_R \frac{1}{r} \sim \delta$$



Fuzzy radial coordinate



- Angular coordinates on layers are properly fuzzy. The radial coordinate is discrete and lattice-like.
- This calls for improvement. Several possible ways how to do this.



[Kovacik, Prekrat, JT work in progress]

$$\mathcal{S}\phi^{(N)} = \frac{\phi^{(N)} + \sum_i \alpha_i (\mathcal{U}^i \phi^{(N-i)} + \mathcal{D}^i \phi^{(N+i)})}{1 + \sum_i \alpha_i}$$

- This procedure simply smears the values of fields $\phi^{(N)}$ over neighboring layers.



[Kovacik, Prekrat, JT work in progress]

$$\mathcal{S}\phi^{(N)} = \frac{\phi^{(N)} + \sum_i \alpha_i (\mathcal{U}^i \phi^{(N-i)} + \mathcal{D}^i \phi^{(N+i)})}{1 + \sum_i \alpha_i}$$

- This procedure simply smears the values of fields $\phi^{(N)}$ over neighboring layers.
- What does this do?



- Functions on the fuzzy sphere are matrices acting on \mathcal{H}

$$\Phi = \sum_{m,n=-s}^s \Phi_{mn} |m\rangle \langle n| .$$

- We can express the matrix Φ in a similar fashion using the coherent states

$$\Phi = \left(\frac{N}{4\pi} \right)^2 \int d^2x d^2y \phi(x,y) |x\rangle \langle y| .$$

- Objects [\[Iso, Kawai, Kitazawa 2000; Steinacker 2016; Steinacker, JT '22\]](#)

$$|x\rangle \langle y| =: \begin{vmatrix} x \\ y \end{vmatrix}$$

form a basis of functions on the fuzzy sphere and we will call them the **string modes**.



FUZZY RADIAL COORDINATE – STRING STATES

- In onion construction, for x and y on the same layer these for matrices $\Phi^{(N)}$.
- For x and y on different layers these naturally fit into the off-diagonal blocks of

$$\Psi = \begin{pmatrix} \Phi^{(1)} & \bullet & & \\ \bullet & \Phi^{(2)} & & \\ & & \ddots & \\ & & & \Phi^{(M)} \end{pmatrix}.$$



FUZZY RADIAL COORDINATE – STRING STATES

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$$\Psi = \begin{pmatrix} \Phi^{(1)} & \bullet & & \\ \bullet & \Phi^{(2)} & & \\ & & \ddots & \\ & & & \Phi^{(M)} \end{pmatrix}.$$

- What is the effect of off-diagonal blocks? What is their correct dynamics? Is this with or instead of the radial derivative?
- Can we define physics of the onion in terms of the whole matrix Ψ ?



- A simplified model of two spheres with the same N

$$\Psi_F = \begin{pmatrix} \phi_1 & A \\ A^\dagger & \phi_2 \end{pmatrix} .$$

- We take A to be identity and get

$$\text{Tr} \left[\frac{1}{2} r \Psi^2 + \Psi^4 \right] = \frac{1}{2} (r + 4) \text{tr}_N (\Phi_1^2) + \text{tr}_N (\Phi_1^4) + \frac{1}{2} (r + 4) \text{tr}_N (\Phi_2^2) + \text{tr}_N (\Phi_2^4) + 4 \text{tr}_N (\Phi_1 \Phi_2)$$

- This is a solvable two matrix model related to Ising model.



- There is a different and more mathematical description of noncommutative spaces.
- Uses notion of spectral triples

$$(\mathcal{A}, \mathcal{D}, \mathcal{H}) .$$

- Construction for fuzzy sphere available, construction of a lattice like set of points available.
[\[Barrett '15\]](#)



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- Construction for fuzzy sphere available, construction of a lattice like set of points available.
[\[Barrett '15\]](#)
- Can we put these two together? Any other (better) version for radial direction? Does this lead to any canonical structures on the onion?



Onion as a spacetime solution in matrix models



ONION AS A SPACETIME SOLUTION IN MATRIX MODELS

- Matrix models formulations of string theory usually have action along the lines

$$S = \frac{1}{g^2} \text{Tr} [-[X_a, X_b][X^a, X^b] + \dots]$$

with matrices X_a describing spacetime degrees of freedom. [\[Steinacker '24\]](#)

- Equations of motion

$$[X_a, [X^a, X_b]] = \dots$$

lead to solutions in form of fuzzy spaces. The simplest case is set of fuzzy spheres of various radii

$$X_a = \begin{pmatrix} L_a^{(N_1)} & & & \\ & L_a^{(N_2)} & & \\ & & \ddots & \\ & & & L_a^{(N_M)} \end{pmatrix}.$$

- Hmmmmmmmm.



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- Hmmmmmmmm.
- What does this tell us about the construction of fuzzy onion? What about the off-diagonal blocks?



Model of dynamical spacetime



MODEL OF DYNAMICAL SPACETIME

- In the construction λ is a constant. Does not need be and in principle we could have $\lambda(r)$ and a deformed onion – curvature.
- This deformation does not need to be constant in time.
- Perhaps also possibility of making λ dependent on the angular direction.



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- This deformation does not need to be constant in time.
- Perhaps also possibility of making λ dependent on the angular direction.
- What kind of dynamics of space can we define? What is the effect on physics happening on the onion?



TOY MODEL OF EXPANDING UNIVERSE

- If we interpret the radial coordinate as time, fuzzy onion is a model of expanding universe with quantized time.
- Each time step one cell of spacetime is created.



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- If we interpret the radial coordinate as time, fuzzy onion is a model of expanding universe with quantized time.
- Each time step one cell of spacetime is created.
- What are the consequences?



Phenomenology of infalling matter



PHENOMENOLOGY OF INFALLING MATTER

- Properties of black holes in quantum spacetimes are different, no singularity present.
- We can analyze collapse of matter creating a black hole by writing the corresponding equations.



- Properties of black holes in quantum spacetimes are different, no singularity present.
- We can analyze collapse of matter creating a black hole by writing the corresponding equations.
- Is the collapse stopped by outward pressure? What is the dissipation mechanism? What is fate of the horizon? Any bounces?
- Comparison with numerical results in loop quantum gravity.
[Modesto '08; Husain, Kelly, Santacruz, Wilson-Ewing '22]



Classical applications



- The cells of the "quantum" space need to arise from fundamental physics.
- Flow of granular materials or heat flow in such materials. [\[Saitou, Bamba, Sugamot '14\]](#)
- Structure of neutron stars.
- Applicable in situations where granularity is due to lack of precise knowledge – atmospheric physics.



Take home message and 2do list



TAKE HOME MESSAGE AND 2DO LIST

- Quantization of gravity seems to lead to quantized spacetimes.
- Fuzzy onion is a three dimensional model of such a space.
- It is rather straightforward to work with so it is a nice toy model / playground to check the consequences of quantum structure.



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Thank you for your attention!



If time permits



FIELD THEORY AS A VECTOR MODEL

- Reformulation in terms of vector \mathcal{C} and operators as $d \times d$ matrices.
- Is this useful for the field theory in any way?
- Action turns out to be

$$\begin{aligned} S &= 4\pi \text{Tr} \, r \left(a \Psi K \Psi + b \Psi^2 + c \Psi^4 \right) = \\ &= \frac{1}{2} \mathcal{C} \cdot \mathbf{P}^{-1} \cdot \mathcal{C} + 4\pi \lambda^3 \sum_{N=1}^M c_N \left[\frac{1}{4N} \left(\mathcal{C}^{(N)} \cdot \mathcal{C}^{(N)} \right)^2 + \frac{1}{8} \left(\mathcal{C}^{(N)} \cdot \mathbf{G}_a^{(N)} \cdot \mathcal{C}^{(N)} \right)^2 \right], \\ \mathbf{P} &= \frac{1}{4\pi \lambda^2} (2arK + 2br)^{-1} \end{aligned}$$

where

$$\mathbf{G}_a^{(N)} = \begin{pmatrix} 0 & (v_a^{(N)})^T \\ v_a^{(N)} & D_a^{(N)} \end{pmatrix}, \quad (D_a^{(N)})_{ij} = 2 \text{tr}_N \left(\left\{ T_i^{(N)}, T_j^{(N)} \right\} T_a^{(N)} \right), \quad (v_a^{(N)})_b = \sqrt{\frac{2}{N}} \delta_{ab}.$$



- What are consequences of writing

$$S = \frac{1}{2} \mathcal{C} \cdot \mathbf{P}^{-1} \cdot \mathcal{C} + 4\pi c \lambda^3 \left[\frac{1}{4} (\mathcal{C} \cdot \mathcal{C})^2 + \frac{1}{8} (\mathcal{C} \cdot \mathbf{G}_A \cdot \mathcal{C})^2 \right]$$

instead of

$$S = \frac{1}{2} \mathcal{C} \cdot \mathbf{P}^{-1} \cdot \mathcal{C} + 4\pi \lambda^3 \sum_{N=1}^M c_N \left[\frac{1}{4N} \left(\mathcal{C}^{(N)} \cdot \mathcal{C}^{(N)} \right)^2 + \frac{1}{8} \left(\mathcal{C}^{(N)} \cdot \mathbf{G}_a^{(N)} \cdot \mathcal{C}^{(N)} \right)^2 \right]?$$

